

Surface Science 417 (1998) 406-414

# CO<sub>2</sub> chemisorption at Ca and CaO surfaces: a study with MIES, UPS(HeI) and XPS

D. Ochs, B. Braun, W. Maus-Friedrichs, V. Kempter \*

Physikalisches Institut der Technischen Universität Clausthal, Leibnizstraße 4, 38678 Clausthal-Zellerfeld, Germany

Received 10 July 1998; accepted for publication 1 September 1998

#### Abstract

Metastable impact electron spectroscopy (MIES) in combination with UPS(HeI) and X-ray photoelectron spectroscopy was applied to study chemisorption of CO<sub>2</sub> on Ca and CaO films at room temperature. Ca surfaces were prepared by evaporating Ca films (ca 20 nm thickness) on a clean Si(111) surface; information on the topological structure of the Ca film was obtained with scanning tunneling microscopy (STM). CaO films were produced by exposing Ca films to O<sub>2</sub> or N<sub>2</sub>O. Ca surfaces exposed to CO<sub>2</sub> become terminated by carbonate complexes ( $CO_3^{2-}$ ) while features typical for (Ca–O) bonds are detected underneath the surface. A closed carbonate layer forms also when a CaO film is exposed to CO<sub>2</sub>. The strong reactivity of the CaO surface against CO<sub>2</sub> is in agreement with recent results of cluster calculations. © 1998 Elsevier Science B.V. All rights reserved.

*Keywords:* Calcium; Calcium oxide; CO<sub>2</sub>; Metastable impact electron spectroscopy (MIES); Photoelectron spectroscopy (UPS(HeI)); Scanning tunneling microscopy (STM)

#### 1. Introduction

The reaction of molecules with oxide surfaces is of great interest for the development of new gas sensors and devices in heterogeneous catalysis. Nevertheless, a fundamental understanding of chemisorption and catalysis is still limited in these cases. So far, the most popular oxide surfaces for gas sensor applications are ZnO and SnO<sub>2</sub> [1]. As far as CaO is concerned, the interaction of SO<sub>2</sub> with CaO(100) has been investigated in some detail with X-ray photoelectron spectroscopy (XPS) under ultrahigh vacuum (UHV) conditions [2,3]. The reactivity of large-area CaO powder with color centers against molecular oxygen was investigated with electron paramagnetic resonance [4]. Cluster calculations modeling the interaction of CO<sub>2</sub> and SO<sub>2</sub> with CaO and MgO were performed by Pacchioni et al. [5,6]. It was found that CO<sub>2</sub> chemisorbs as carbonate molecular ion  $(CO_3^{3-})$  at oxygen sites on both surfaces. The formation of a monodentate carbonate species was predicted. While CO<sub>2</sub> chemisorbs on MgO at low-coordinated oxygen sites only, the carbonate formation occurs also at regular oxygen sites on CaO. The prediction made for MgO is supported by our previous measurements [7,8].

In this article we study the chemisorption of  $CO_2$  on CaO surfaces applying essentially the same strategy as in our previous studies concerned with the chemisorption of  $CO_2$  on MgO [7,8]. At first, we describe ways to produce CaO films starting with a Ca surface. Then, we study the interaction

<sup>\*</sup> Corresponding author. Fax: +49 5323 723600; e-mail: volker.kempter@tu-clausthal.de

<sup>0039-6028/98/\$ –</sup> see front matter  $\odot$  1998 Elsevier Science B.V. All rights reserved. PII: S0039-6028(98)00721-3

of the Ca and CaO films with CO<sub>2</sub>. We characterize the electronic structure of Ca and CaO surfaces, before and after exposure to CO<sub>2</sub>, by combining metastable impact electron spectroscopy (MIES), photoelectron spectroscopy UPS(HeI) and XPS. The identification of the CO<sub>2</sub> induced features is accomplished from a comparison with the spectral features developed when CO<sub>2</sub> interacts with Ca surfaces. This task is facilitated by the information already available on the CO<sub>2</sub> interaction with other metal surfaces [9,10]. Information on the topology of the Ca surface comes from scanning tunneling microscopy (STM).

### 2. Experimental

The experiment was carried out in an UHV system (base pressure  $< 2 \times 10^{-10}$  Torr) consisting of a chamber for electron spectroscopies and another one for sample treatment [7,11,12]. The former one is equipped with facilities for XPS, low energy electron diffraction (LEED) and Auger electron spectroscopy (AES) and a cold-cathode gas discharge source for the production of metastable He\* $(2^{3}S/2^{1}S)$  (E\*=19.8/20.6 eV) atoms with thermal kinetic energies as a source for MIES and HeI photons ( $E^* = 21.2 \text{ eV}$ ) as a source for UPS [11] (and references cited therein). The triplet-tosinglet ratio has been measured by He\*-Ar impact as 7:1. Metastable and photon contributions within the beam are separated by means of a time-offlight technique using a mechanical chopper in combination with a double counter system allowing to measure MIES and UPS quasi-simultaneously. All electron energy spectra were acquired with the incident photon/metastable beams under 45° with respect to surface normal in the constant pass energy mode using a hemispherical analyzer with an acceptance angle of  $8^{\circ}$  (FWHM); the binding energies are referenced to the Fermi level  $E_{\rm F}$ . By biasing the sample with respect to the analyzer (50 eV) the variation of the surface work function with exposure can be determined directly from the high binding energy cutoff of the electron spectra. For binding energies larger than ca 10 eV both the UPS and the MIES spectra appear to be influenced by secondary electrons; this part of the spectra will not be discussed in the following.

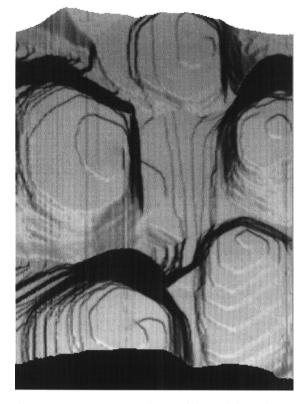
The STM apparatus consists of a commercial Perkin–Elmer chamber (base pressure  $<1\times$ 10<sup>-10</sup> Torr) equipped with an Omicron Micro-STM [13]. It is furthermore equipped with a cylindrical mirror analyzer (CMA) to perform AES. The target is cleaned in a preparation stage in order to prevent an eventual contamination of the STM. The STM images were taken in the constant current topography (CCT) mode. The bias voltage given with the images corresponds to the target potential. The STM tip is fixed at zero potential. The calibration of the scanner z-coordinate is known from a previous study [13]. The n-doped Si(111) substrate was prepared from commercial wafers. They were cleaned by electron bombardment, heating to ca 1200°C for several minutes with subsequent cooling down to a temperature of ca 600°C within 10 min. After this procedure the silicon substrates showed the Si(111)- $(7 \times 7)$  reconstruction both in LEED and STM. Ca layers were produced by evaporating calcium from a commercial Knudsen cell in the MIES apparatus and a home-built evaporator in the STM chamber. During the evaporation of Ca the pressure never rose beyond  $3 \times 10^{-9}$  mbar. All measurements are performed at room temperature.

# 3. Results

#### 3.1. Characterization of Ca and CaO films

Fig. 1 shows the STM image of the Si(111) surface covered by a closed Ca film (total thickness > 20 nm). The surface is terminated by threedimensional islands. Their lateral size is between 50 and 100 nm; their height is ca 5 nm, typically. The step height between each two terraces of the hexagonal islands is  $2.64 \pm 0.2$  Å, which is close to half of the Ca(fcc) lattice constant (2.79 Å [14]). Furthermore, most islands are aligned relative to each other. Although atomic resolution could not be achieved, our findings appear to suggest that the islands consist of well ordered Ca atoms.

The MIES spectra of Ca films on Si(111) exposed to oxygen are presented in Fig. 2. The



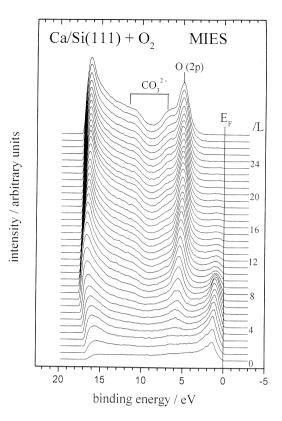


Fig. 1. 300 nm  $\times$  350 nm STM image of the Ca film on Si(111); max. height difference, 10 nm; bias voltage, -2 eV; tunneling current, 1.2 nA.

bottom spectrum applies to the clean Ca film and shows a well-developed Fermi edge. This is characteristic for spectra caused by Auger deexcitation. Hereby, an electron from the uppermost occupied Ca band fills the He 1s hole of the probe atom while the He 2s electron is emitted, carrying the excess energy [15,16]. The work function of the Ca film, derived from the low energy onset of the bottom MIES spectrum, is  $3.2\pm0.2$  eV; the value for Ca(100) is 2.8 eV [15].

The presence of oxygen at the surface manifests itself in two respects:

- (1) in the development of a new peak at ca  $E_{\rm B} = 5.5 \, {\rm eV}$ ; and
- (2) in the spectral changes and, finally, the disappearance of the structure close to  $E_{\rm F}$  which is representative for the metallic part of the surface density of states (SDOS).

By comparison with the corresponding structure

Fig. 2. MIES spectra of Ca film (thickness > 20 nm) on Si(111) exposed to O<sub>2</sub>.

observed during the oxygen exposure of Mg films [8,12] we attribute the feature seen at ca  $E_{\rm B}$  = 5.5 eV to the ionization of O 2p of oxygen atoms chemisorbed at the Ca surface.

The spectral changes upon oxygen exposure observed in the metallic part of the SDOS can probably be understood as follows: in the initial stage of the exposure the s-part of the SDOS becomes preferentially involved in the bonding of the oxygen, and its contribution to the Ca charge density is strongly reduced. Therefore, MIES samples the higher angular moment components of the metallic part of the surface charge density which are less involved in (Ca–O) bonds. In particular, the p- and e<sub>g</sub>-part of the DOS is strongly peaked at  $E_{\rm F}$  [17], and may therefore be responsible for the narrowing of the peak seen near  $E_{\rm F}$ .

The formation of a  $He^{-*}(1s2s^2)$  species and its decay by autodetachment offers a second possi-

bility to explain the observed increase in intensity near  $E_{\rm F}$  [18]. However, the work function needs to decrease to values below ca 2.2 eV upon exposure to O<sub>2</sub> or CO<sub>2</sub> for efficient formation of He<sup>-\*</sup> [18]. This is not the case for Ca exposed to either O<sub>2</sub> or CO<sub>2</sub>, and for this reason we expect at most a small contribution of the autodetachment process to the spectra of Figs. 2 and 5.

The work function decreases during the initial stage of the exposure and reaches its minimum ca 10 L (1 L =  $10^{-6}$  Torr s). Around 20 L the metallic part of the SDOS has virtually disappeared. The intensity of the O 2p feature saturates ca 16 L, and the work function has adopted its final value of ca 3 eV.

The energy difference between the top of the valence band and the Fermi level  $E_{\rm F}$  gives the lower part of the bandgap of the insulating layer. For defect free bulk insulators the Fermi level

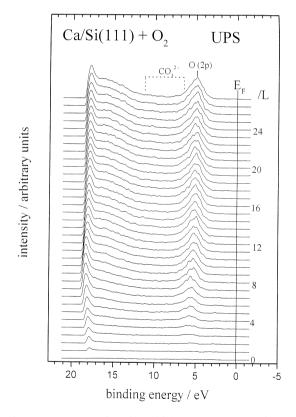


Fig. 3. UPS spectra of Ca film (thickness > 20 nm) on Si(111) exposed to O<sub>2</sub>.

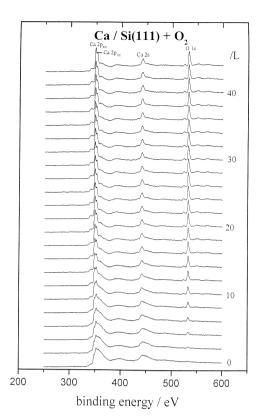


Fig. 4. XPS spectra of Ca film (thickness >20 nm) on Si(111) exposed to O<sub>2</sub>.

would be located in the middle of the band gap. Under this assumption a bandgap width of  $7\pm0.2$  eV is derived from the top spectra of Figs. 2 and 3 (6.9 eV for the CaO bulk [1]).

The distance from  $E_{\rm F}$  to the vacuum level, the work function of the insulating film, is taken from the onset of the spectra at large binding energies. Thus, the distance from the top of the valence band to the vacuum level, as derived from both the MIES and UPS results, is 6.2 eV.

Two additional features, labeled  $CO_3^{2-}$ , appear at  $E_B=7$ ; 11.8 eV which can be attributed to a small surface contamination by carbonate ions (see Section 3.2). We cannot exclude that the presence of an OH-species at the surface also contributes to these features [19].

The UPS results of Fig. 3 are rather insensitive to the metallic part of the DOS, but they confirm the presence of chemisorbed oxygen by the occur-

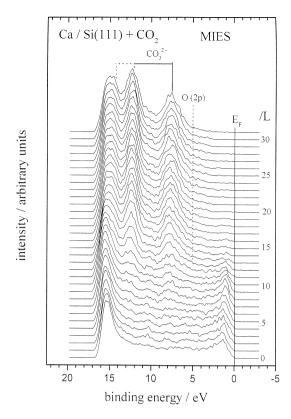


Fig. 5. MIES spectra of Ca film (thickness >20 nm) on Si(111) exposed to CO<sub>2</sub>.

rence of the peak at  $E_B = 5.5$  eV that we attribute to the ionization of O 2p. No features develop in UPS at the positions where carbonate features are expected (see dashed lines) supporting the assumption of the surface origin of the additional species seen with MIES.

Fig. 4 shows the XPS spectra collected for oxygen exposures up to 44 L. The spectrum of the clean Ca film displays two broad structures at  $E_B = 360$ ; 450 eV which can be assigned to the ionization of the Ca 2p and Ca 2s orbitals. Both structures narrow considerably upon oxygen exposure: the Ca 2p peak develops into a double-peak structure at  $E_B = 348.6$ ; 352.0 eV which is caused by the ionization of the two fine structure components Ca  $2p_{3/2}$  and Ca  $2p_{1/2}$ . The Ca 2s structure develops into a peak centered at 441 eV. A new peak appears at  $E_B = 532.2$  eV, due to ionization

of O 1s, displaying a linear rise in intensity up to ca 20 L.

We propose the following interpretation to the oxygen-induced changes:

- The broad structures, assigned to the ionization of the Ca 2p and Ca 2s orbitals, are caused by "shake-up" processes into empty metallic states still available in the initial stage of the oxygen exposure. These processes produce shoulders towards higher binding energies with respect to the main peaks [20]. The states responsible for the "shake-up" processes are not available anymore after the bandgap formation has occurred in the near-surface region.
- The O 1s feature at 532.2 eV saturates at ca 22 L while the O 2p signal in UPS reaches already saturation at ca 18 L [21]. This, together with the fact that no metallic features can be seen beyond ca 20 L, suggest that an insulating film has formed.
- XPS spectra obtained by us from a CaO(100) single crystal (not shown here) are practically identical to the saturation spectra presented in Fig. 4, both as far as the position of the peaks and their relative intensities are concerned [21]. This leads us to the conclusion that indeed a stoichiometric CaO film (thickness >10 nm) has formed.

We can compare our MIES and UPS results to ab initio results obtained by applying the full potential photoemission theory to CaO(100) [22]. The agreement with the full potential angleresolved valence band spectra in the  $\Gamma$ XWK emission plane for unpolarized HeII radiation is rather good, in particular as far as the width and shape of the valence band emission are concerned. This suggests that final state effects may be neglected and that our MIES (UPS) results do mainly represent the SDOS (DOS) of the valence band states, respectively. When Ca films on Si(111) are exposed to N<sub>2</sub>O molecules, the oxygen-induced changes and the equilibrium spectra are, within the accuracy of our results, identical to those presented in Figs. 2-4 although a larger exposure (ca 100 L) is required to reach this stage [21,23]. This finding is compatible with the assumption that the N<sub>2</sub>O molecules dissociate when interacting with the surface. While the N<sub>2</sub> fragments leave the surface, the oxygen atoms form (Ca–O) bonds. Ultimately, the film turns into a CaO layer, just as for  $O_2$  exposure. Again a small surface contamination with carbonate ions is detected with MIES.

Equilibrium spectra identical to those obtained with  $O_2$  and  $N_2O$  can also be obtained from the evaporation of Ca in an oxygen ambient atmosphere [21,23]. The coincidence of the equilibrium spectra from three different methods is further evidence that indeed a CaO film is produced in all three cases.

## 3.2. Ca films exposed to $CO_2$

The MIES and UPS(HeI) results obtained during the CO<sub>2</sub> exposure of a Ca film (>20 nm thick) on Si(111) are presented in Figs. 5 and 6; a small oxygen contamination is seen in the UPS, but not in the MIES spectra. The exposure depen-

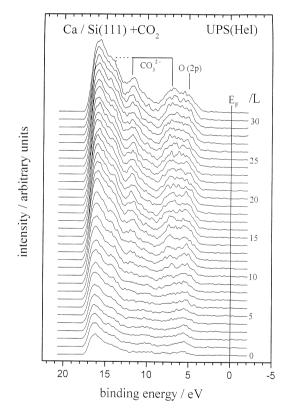


Fig. 6. UPS spectra of Ca film (thickness > 20 nm) on Si(111) exposed to CO<sub>2</sub>.

dence differs considerably from that for  $O_2$  and N<sub>2</sub>O. Moreover, the MIES and UPS results display considerable differences among each other. Peaks, labeled  $CO_3^{2-}$ , appear in the MIES spectra at  $E_{\rm B} = 7.6$ ; 12.2 eV, and possibly at 14.5 eV. They are also seen with UPS where they are superimposed by the feature O 2p. The identification of the spectral features follows from photoelectron spectroscopy results for CaCO<sub>3</sub> bulk samples [24] and from our previous work dealing with the CO<sub>2</sub> interaction with Mg and MgO surfaces where the same features were seen [7,8]. They are due to the ionization of a carbonate  $(CO_3^{2-})$  species, and can be assigned to the MOs  $(1a'_2; 1e'; 4e')$  (7.6 eV),  $(3e';1a'_2)$  (12.2 eV) and  $(4a'_1)$  (14.5 eV) [24]. The emission from  $CO_3^{2-}$  increases up to ca 20 L and, thereafter, remains almost unchanged.

The UPS spectra are more complicated: besides the emission from  $CO_3^{2-}$  which is also seen, the structure labeled O 2p is present. By comparison with the O<sub>2</sub> and N<sub>2</sub>O results we assign it to the ionization of O 2p from atomic oxygen species involved in (Ca–O) bonds. It is remarkable that O 2p is seen in UPS before the carbonate emission becomes visible in MIES.

The XPS spectra of Fig. 7 can be analyzed along the same lines as those of Fig. 4. The metallic character of the surface has disappeared above ca 32 L as judged from the absence of the shake up features. The O 1s peak is split into two components at 532 and 536 eV which we attribute, by comparison with Fig. 4, to oxygen in (Ca–O) bonds (532 eV) and in carbonate complexes (536 eV); the exposure dependence of the peak intensities is more complicated than in Fig. 4.

For reviews of the surface chemistry of  $CO_2$  the reader is referred to Refs [9,10]. Our spectroscopic results obtained with MIES, UPS and XPS show that, as in the case of Mg [9,10], the  $CO_2$  adsorption is dissociative. The presence of oxygen atoms at the surface promotes the bonding of  $CO_2$  as carbonate species. The MIES data, in particular, show that the carbonate and oxide species have rather different distributions in the near-surface regions with the carbonate being confined to the topmost layer and the oxide present as an underlayer. The present results are compatible with the following scenario for the interaction of  $CO_2$ 

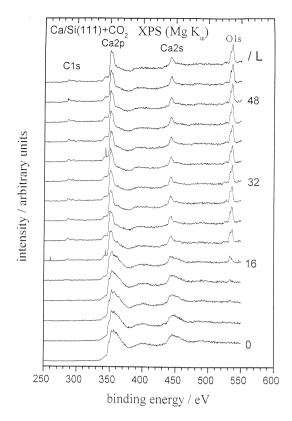


Fig. 7. XPS spectra of Ca film (thickness >20 nm) on Si(111) exposed to CO<sub>2</sub>.

with Ca surfaces first proposed for the interaction of  $CO_2$  with Mg surfaces [25] (see also our previous results for Mg [7,8]):

- (1) the CO<sub>2</sub> molecules dissociate on the surface; in this way (Ca–O) bonds are formed; and
- (2) additional  $CO_2$  can react with surface oxygen which leads to the  $CO_3^{2-}$  termination of the surface.

It is remarkable that, despite the carbonate termination of the surface, a CaO film with >10 nm thickness forms. Additional spectroscopic studies at lower surface temperatures would be desirable for clarification of the details of the reaction dynamics.

# 3.3. CaO exposed to $CO_2$

Fig. 8 shows the MIES results for  $CO_2$  exposure of CaO films on Si(111). We start with a CaO

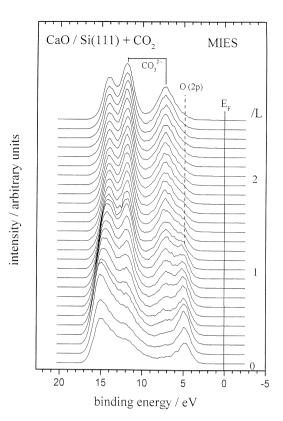


Fig. 8. MIES spectra of CaO film (thickness >10 nm) on Si(111) exposed to CO<sub>2</sub>.

film prepared along the lines described in Section 3.2. The film thickness is ca 10 nm. The MIES spectra of the "as prepared" CaO surface display the density of states in the uppermost filled valence band discussed in the previous sections. After CO<sub>2</sub> exposure the two features labeled  $CO_3^{2-}$  emerge at  $E_B = 7.2$ ; 12.0 eV. They are due to the ionization of a surface carbonate species (see Section 3.2), and are also seen, although only weakly, when exposing MgO surfaces to  $CO_2[7,8]$ . After an exposure of 2 L the two features have developed into strong peaks, and very little contribution from O 2p ionization can be seen in the spectra. In saturation the spectra become very similar to those obtained when exposing the Ca film to oxygen (Section 3.1).

In Fig. 9 the corresponding UPS measurements are shown. Structure  $CO_3^{2-}$  superimposes the O(2p) peak located at ca 5 eV. Under saturation

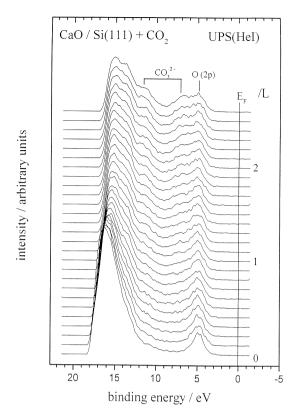


Fig. 9. UPS spectra of CaO film (thickness >10 nm) on Si(111) exposed to CO<sub>2</sub>.

conditions, the structures O 2p and  $CO_3^{2-}$  are of comparable intensity. The spectrum becomes then, as for MIES, very similar to that obtained from exposing the Ca film to oxygen (Section 3.1).

Obviously, a closed  $CO_3^{2-}$  layer forms on top of the CaO surface. For MgO surfaces exposed to  $CO_2$  the carbonate features at ca  $E_B = 7.6$ ; 12.2 eV appear too, but with much less intensity [7,8]; in saturation their intensity is about two orders of magnitude smaller.

The chemisorption of  $CO_2$  on MgO and CaO was studied theoretically by performing cluster calculations [5,6]; it is found that on MgO  $CO_2$  can chemisorb at fourfold-coordinated  $O_{4c}^{2-}$  step sites. It forms a chemisorbed complex whose geometry resembles that of a monodentate carbonate  $CO_3^{2-}$  species. The (O–C–O) unit is considerably distorted from the gas phase linear structure: the (C–O) distance is stretched by 0.06 Å, and the

(O–C–O) angle is reduced to ca 130°. The interaction involves charge transfer from the uppermost filled O 2p states at surface ions into the vacant  $\pi_u^*$  MO of CO<sub>2</sub>. This charge transfer occurs more easily at a low coordinated defect site,  $O_{4c}^{2-}$ , than at regular  $O_{5c}^{2-}$  sites: at fourfold-coordinated  $O_{4c}^{2-}$ step sites the electrons are less strongly bound, and charge can flow to the adsorbed CO<sub>2</sub>.

A dramatic difference was predicted for the reactivity of CaO against CO<sub>2</sub>: strong adsorption with a binding energy of ca 1 eV takes place at regular  $O_{5c}^{2-}$  sites of the surface. The surface complex is not very different from that formed in the case of MgO at low-coordinated  $O_{4c}^{2-}$  defect sites. According to Refs [5,6] the strongly different reactivity of the two oxide surfaces can be ascribed to the different Madelung potentials of the two surfaces. This has two important consequences:

- (1) the electron cloud of the  $O_2^-$  anions in CaO is spatially more diffuse than in MgO and can overlap more efficiently with the unoccupied  $\pi_u^-$  of the incoming CO<sub>2</sub> molecules; and
- (2) the HOMO donor orbital  $O_2^-$  lies at ca 11 eV in MgO, but at ca 8.5 eV in CaO.

Thus, charge can be donated to the adsorbing  $CO_2$  molecule at a lower cost in CaO than in MgO. As a consequence, the Madelung potential at a regular  $O_{5c}^{2-}$  site of the CaO surface is intermediate between that of a step and a corner defect site of MgO.

The numerical results find support by our results for the CO<sub>2</sub> chemisorption on MgO [7,8] and the present work on CaO: for MgO a saturation carbonate coverage of the order of only 1% of the number of available oxygen sites is obtained at room temperature. On the contrary, the number of carbonate complexes is about two orders of magnitude larger for CaO.

No chemisorbed species, as f.i.  $CO_2^{\delta^-}$  and  $O_2^{\delta^-}$ , could be detected with MIES when CaO films, prepared as described in Section 3.2, were exposed to CO and O<sub>2</sub>, respectively [23].

#### 4. Summary

Ca films were exposed to  $O_2$  and  $N_2O$  molecules at room temperature in order to produce a CaO film. Information on the topological structure of the Ca film was obtained with STM. The electronic structure of the Ca film and the structural changes taking place during gas exposure were studied with the electron spectroscopic techniques MIES, UPS(HeI) and XPS. The results indicate that a stoichiometric CaO film, >10 nm thick, can be produced in this way.

Ca exposed to CO<sub>2</sub> at room temperature develops spectral features which arise from the ionization of the  $(1a'_2; 1e'; 4e')$ ,  $(3e'; 1a'_2)$  and  $(4a'_1)$ molecular orbitals of the carbonate molecule ion  $(CO_3^{2-})$ . UPS shows the same features, but superimposed by a peak which is assigned to the ionization of O 2p. The results support the following sequence of processes for the interaction of CO<sub>2</sub> with Ca:

- the CO<sub>2</sub> dissociates upon interacting with the surface whereby (Ca–O) bonds are formed; and
- (2) additional  $CO_2$  reacts with surface oxygen which leads to  $CO_2$  chemisorption in form of  $CO_3^{2-}$ .

This mechanism is the same as proposed for the  $CO_2$  interaction with Mg surfaces [7,8,25].

A closed carbonate layer forms when a CaO film is exposed to  $CO_2$  at room temperature. This is in contrast to MgO where  $CO_2$  chemisorbs at low-coordinated oxygen sites only [7,8]. According to Refs [5,6] the strongly different reactivity of the two surfaces can be attributed to the different Madelung potentials with the consequence that charge donation to the chemisorbing  $CO_2$  can occur at a lower cost in CaO than in MgO.

#### Acknowledgements

Financial assistance came from the DAAD/BC project ARC 313-ARC-XI-97/51. The authors acknowledge fruitful discussions with A. Shluger and L. Kantorovich (University College London).

## References

- V.E. Henrich, P.A. Cox, The Surface Science of Metal Oxides, Cambridge University Press, Cambridge, 1994.
- [2] Y.C. Lee, P.A. Montano, J.M. Cook, Surf. Sci. 143 (1984) 423.
- [3] C.D. Stinespring, J.M. Cook, J. Electron Spectrosc. Relat. Phenom. 32 (1983) 113.
- [4] M. Chiesa, M.C. Paganini, E. Giamello, Langmuir 13 (1997) 5306.
- [5] G. Pacchioni, Surf. Sci. 281 (1993) 207.
- [6] G. Pacchioni, J.M. Ricart, F. Illas, J. Am. Chem. Soc. 116 (1994) 10152.
- [7] D. Ochs, M. Brause, B. Braun, W. Maus-Friedrichs, V. Kempter, Surf. Sci. 397 (1998) 184.
- [8] D. Ochs, M. Brause, W. Maus-Friedrichs, V. Kempter, J. Electron Spectrosc. Relat. Phenom. 88 (1998) 757.
- [9] F. Solymosi, J. Molec. Cat. 65 (1991) 337.
- [10] H.-J. Freund, M.W. Roberts, Surf. Sci. Rep. 25 (1996)
  225; J. Electron. Spectrosc. Relat. Phenom. 18 (1980) 23.
- [11] W. Maus-Friedrichs, M. Wehrhahn, S. Dieckhoff, V. Kempter, Surf. Sci. 237 (1990) 257.
- [12] D. Ochs, W. Maus-Friedrichs, M. Brause, J. Günster, V. Kempter, V. Puchin, A. Shluger, L. Kantorovich, Surf. Sci. 365 (1996) 557.
- [13] J. Hauschild, H.J. Elmers, U. Gradmann, Phys. Rev. B 57 (1998) R677.
- [14] R.W.G. Wyckoff, Crystal Structures, Wiley, 1965.
- [15] G. Ertl, J. Küppers, Low Energy Electrons and Surface Chemistry, VCH, Weinheim, Berlin, 1985.
- [16] V. Kempter, Comments on Atomic and Molecular Processes 34 (1998) 11.
- [17] D.A. Papaconstantopoulos, Handbook of the Band Structure of Elemental Solids, Plenum Press, New York, 1986.
- [18] J. Günster, Th. Mayer, V. Kempter, Surf. Sci. 359 (1996) 155.
- [19] J. Günster, G. Liu, V. Kempter, D.W. Goodman, J. Vac. Sci. Technol. A 16 (1998) 996.
- [20] D. Briggs, M.P. Seah (Eds.), Practical Surface Analysis by Auger and X-ray Photoelectron Spectroscopy, Wiley, Chichester, 1983.
- [21] B. Braun, Diploma Thesis, TU Clausthal, unpublished.
- [22] M. Grass, J. Braun, G. Borstel, Surf. Sci. 334 (1995) 215.
- [23] D. Ochs, Doctoral Thesis, TU Clausthal, 1998.
- [24] E. Tegeler, N. Kosuch, G. Wiech, A. Faessler, J. Electron Spectrosc. Relat. Phenom. 18 (1980) 23.
- [25] S. Campbell, P. Hollins, E. McCash, M.W. Roberts, J. Electron Spectrosc. Relat. Phenom. 145 (1986) 145.